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Envelopes of slant lines in the hyperbolic plane

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Abstract. In this paper we consider envelopes of families of equidistant curves and horocycles in the hyperbolic plane. As a special case, we consider a kind of evolutes as the envelope of normal equidistant families of a curve. The hyperbolic evolute of a curve is a special case. Moreover, a new notion of horocyclic evolutes of curves is induced. We investigate the singularities of such envelopes and introduce new invariants in the Lie algebra of the Lorentz group.

Keywords: slant geometry, Hyperbolic plane, horocycles, equidistant curves

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1 Introduction

We consider the Poincaré disk model D of the hyperbolic plane which is conformally equivalent to the Euclidean plane, so that a circle or a line in the Poincaré disk is also a circle or a line in the Euclidean plane. A geodesic in the Poincaré disk is a Euclidean circle or a line which is perpendicular to the ideal boundary (i.e., the unit circle). If we adopt geodesics as lines in the Poincaré disk, we have the model of the hyperbolic geometry. A *horocycle* is an Euclidean circle which is tangent to the ideal boundary. If we adopt horocycles as lines, we call this geometry a *horocyclic geometry* (a *horospherical geometry* for the higher dimensional case) [4, 6, 7, 8, 9]. We also have another kind of curves with the properties similar to those of Euclidean lines. A curve in the Poincaré disk is called an *equidistant curve* if it is a Euclidean circle or a Euclidean line whose intersection with the ideal boundary consists of two points. We define an equidistant curve depends on $\phi \in [0, \pi/2]$ whose angles with the ideal boundary

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at the intersection points are ϕ (cf., [10]). A geodesic is the special case with $\phi = \pi/2$ and a horocycle is the case with $\phi = 0$. Therefore, a geodesic is called a *vertical pseudo-line* and a horocycle a *horizontal pseudo-line*. For $\phi \in (0, \pi/2]$, the corresponding pseudo-line is an equidistant curve, which we call a ϕ -slant *pseudo-line*. If we consider a ϕ -slant pseudo-line as a line, we call this geometry a *slant geometry* (cf., [1]).

In this paper we consider envelopes of families of ϕ -slant pseudo-lines in the general setting. We investigate the singularities of such envelopes. Throughout the remainder of the paper, we adopt the Lorentz-Minkowski space model of the hyperbolic plane. For a 3×3 -matrix A, we say that A is a member of the Lorentz group $SO_0(1,2)$ if det A > 0 and the induced linear mapping preserves the Lorentz-Minkowski scalar product. The Lorentz group $SO_0(1,2)$ canonically acts on the hyperbolic plane. It is well known that this action is transitive, so that the hyperbolic space is canonically identified with the homogeneous space $SO_0(1,2)/SO(2)$. It follows that any point of the hyperbolic space can be identified with a matrix $A \in SO_0(1,2)$ (cf., §3). Therefore, a one parameter family of ϕ -slant pseudo-lines can be parametrized by using a curve in $SO_0(1,2)$ (cf., §3 and 4). Then we apply the theory of unfoldings of function germs (cf., [2]) and obtain a classification of singularities of the envelopes of the families of ϕ -slant pseudo-lines (cf., Theorem 5.6). The singularities of the envelopes are characterized by using invariants represented by the elements of Lie algebra $\mathfrak{so}(1,2)$ of $SO_0(1,2)$. In §6 we introduce the notion of ϕ -slant evolutes of unit speed curves in the hyperbolic plane. If $\phi = \pi/2$, then the ϕ -slant evolute is a hyperbolic evolutes defined in [5]. Moreover, if $\phi = 0$, then the ϕ -slant evolute is called a *horocyclic evolute*. It means that the ϕ -slant evolutes depending on ϕ connects the hyperbolic evolute and the horocyclic evolute of the curve in the hyperbolic plane.

In [3] families of equal-angle envelopes in the Euclidean plane is investigated.

2 Basic concepts

We now present basic notions on Lorentz-Minkowski 3-space. Let $\mathbb{R}^3 = \{(x_0, x_1, x_2) | x_i \in \mathbb{R}, i = 0, 1, 2\}$ be a 3-dimensional vector space. For any vectors $\boldsymbol{x} = (x_0, x_1, x_2), \boldsymbol{y} = (y_0, y_1, y_2) \in \mathbb{R}^3$, the *pseudo scalar product* (or, the Lorentz-Minkoski scalar product) of \boldsymbol{x} and \boldsymbol{y} is defined by $\langle \boldsymbol{x}, \boldsymbol{y} \rangle = -x_0 y_0 + x_1 y_1 + x_2 y_2$. The space $(\mathbb{R}^3, \langle, \rangle)$ is called Lorentz-Minkowski 3-space which is denoted by \mathbb{R}^3_1 . We assume that \mathbb{R}^3_1 is time-oriented and choose $\boldsymbol{e}_0 = (1, 0, 0)$ as the future timelike vector.

We say that a non-zero vector \boldsymbol{x} in \mathbb{R}^3_1 is *spacelike*, *lightlike* or *timelike* if $\langle \boldsymbol{x}, \boldsymbol{x} \rangle > 0 = 0$ or < 0 respectively. The norm of the vector $\boldsymbol{x} \in \mathbb{R}^3_1$ is defined

by $\|\boldsymbol{x}\| = \sqrt{|\langle \boldsymbol{x}, \boldsymbol{x} \rangle|}$. Given a non-zero vector $\boldsymbol{n} \in \mathbb{R}^3_1$ and a real number c, the plane with pseudo normal \boldsymbol{n} is given by

$$P(oldsymbol{n},c) = \{oldsymbol{x} \in \mathbb{R}^3_1 | \langle oldsymbol{x},oldsymbol{n}
angle = c \}.$$

We say that $P(\mathbf{n}, c)$ is *spacelike*, *timelike* or *lightlike* if \mathbf{n} is timelike, spacelike or lightlike respectively.

For any vectors $\boldsymbol{x} = (x_0, x_1, x_2)$, $\boldsymbol{y} = (y_0, y_1, y_2) \in \mathbb{R}^3_1$, pseudo exterior product of \boldsymbol{x} and \boldsymbol{y} is defined to be

$$oldsymbol{x} \wedge oldsymbol{y} = egin{bmatrix} -oldsymbol{e}_0 & oldsymbol{e}_1 & oldsymbol{e}_2\ x_0 & x_1 & x_2\ y_0 & y_1 & y_2 \end{bmatrix} = (-(x_1y_2 - x_2y_1), x_2y_0 - x_0y_2, x_0y_1 - x_1y_0),$$

where $\{e_0, e_1, e_2\}$ is the canonical basis of \mathbb{R}^3_1 . We also define *Hyperbolic plane* by

$$H^{2}_{+}(-1) = \{ \boldsymbol{x} \in \mathbb{R}^{3}_{1} | \langle \boldsymbol{x}, \boldsymbol{x} \rangle = -1, x_{0} \ge 1 \},\$$

de Sitter 2-space by

$$S_1^2 = \{oldsymbol{x} \in \mathbb{R}^3_1 | \langle oldsymbol{x}, oldsymbol{x}
angle = 1 \}$$

and the (open) lightcone at the origin by

$$LC^* = \{ \boldsymbol{x} = (x_0, x_1, x_2) \in \mathbb{R}^3_1 | x_0 \neq 0, \langle \boldsymbol{x}, \boldsymbol{x} \rangle = 0 \}.$$

We remark that $H^2_+(-1)$ is a Riemannian manifold if we consider the induced metric from \mathbb{R}^3_1 .

We now consider the plane defined by $\mathbb{R}_0^2 = \{(x_0, x_1, x_2) \in \mathbb{R}_1^3 \mid x_0 = 0\}$. Since $\langle, \rangle|_{\mathbb{R}_0^2}$ is the canonical Euclidean scalar product, we call it *Euclidean plane*. We adopt coordinates (x_1, x_2) of \mathbb{R}_0^2 instead of $(0, x_1, x_2)$. On Euclidean plane \mathbb{R}_0^2 , we have the *Poincaré disc model* of the hyperbolic plane. We consider a unit open disc $D = \{ \boldsymbol{x} \in \mathbb{R}_0^2 \mid ||\boldsymbol{x}|| < 1 \}$ and consider a Riemannian metric

$$ds^{2} = \frac{4(dx_{1}^{2} + dx_{2}^{2})}{1 - x_{1}^{2} - x_{2}^{2}}.$$

Define a mapping $\Psi : H^2_+ \longrightarrow D$ by

$$\Psi(x_0, x_1, x_2) = \left(\frac{x_1}{x_0 + 1}, \frac{x_2}{x_0 + 1}\right).$$

It is known that Ψ is an isometry. Moreover, the Poinaré disc model is conformally equivalent to the Euclidean plane.

3 Pseudo-lines in the hyperbolic plane

We consider a curve defined by the intersection of the hyperbolic plane with a plane in Lorentz-Minkowski 3-space, which is called a *pseudo-circle* if it is nonempty. The image of a pseudo-circle by the isometry Ψ is a part of a Euclidean circle in the Poincaré disc D. Let $P(\boldsymbol{n}, c)$ be a plane with a unit pseudo-normal \boldsymbol{n} . We call $H^2_+(-1) \cap P(\boldsymbol{n}, c)$ a circle, an equidistant curve and a horocyle if \boldsymbol{n} is timelike, spacelike or lightlike respectively. Moreover, if \boldsymbol{n} is spacelike and c = 0, then we call it a hyperbolic line (or, a geodesic). We remark that circles are compact and other pseudo-circles are non-compact. Therefore, equidistant curves or horocycles are called *pseudo-lines*.

We now consider a hyperbolic line

$$HL(oldsymbol{n}) = \{oldsymbol{x} \in H^2_+(-1) \mid \langle oldsymbol{x}, oldsymbol{n}
angle = 0\}$$

and a horocycle

$$HC(\boldsymbol{\ell},-1) = \{\boldsymbol{x} \in H^2_+(-1) \mid \langle \boldsymbol{x}, \boldsymbol{\ell} \rangle = -1\},\$$

where $\boldsymbol{\ell}$ is a lightlike vector. In general, a horocycle is defined by $\langle \boldsymbol{x}, \boldsymbol{\ell} \rangle = c$ for a lightlike vector $\boldsymbol{\ell}$ and $c \neq 0$. However, if we choose $-\boldsymbol{\ell}/c$ instead of $\boldsymbol{\ell}$, then we have the above equation. We now consider parametrizations of a horocycle and a hyperbolic line respectively. For any $\boldsymbol{a}_0 \in HC(\boldsymbol{\ell}, -1)$, let \boldsymbol{a}_1 be a unit tangent vector of $HC(\boldsymbol{\ell}, -1)$ at \boldsymbol{a}_0 , so that $\langle \boldsymbol{a}_1, \boldsymbol{\ell} \rangle = 0$. We define $\boldsymbol{a}_2 = \boldsymbol{a}_0 \wedge \boldsymbol{a}_1$. Then we have a pseudo orthonormal basis $\{\boldsymbol{a}_0, \boldsymbol{a}_1, \boldsymbol{a}_2\}$ of \mathbb{R}_1^3 such that $\langle \boldsymbol{a}_0, \boldsymbol{a}_0 \rangle = -1$. We remark that \boldsymbol{a}_0 is timelike and $\boldsymbol{a}_1, \boldsymbol{a}_2$ are spacelike. Since $\langle \boldsymbol{\ell} - \boldsymbol{a}_0, \boldsymbol{a}_0 \rangle =$ $\langle \boldsymbol{\ell}, \boldsymbol{a}_1 \rangle = 0$, we have $\pm \boldsymbol{a}_2 = \boldsymbol{\ell} - \boldsymbol{a}_0$. We choose the direction of \boldsymbol{a}_1 such that $\boldsymbol{a}_2 = \boldsymbol{\ell} - \boldsymbol{a}_0$. It follows that $A = ({}^t\boldsymbol{a}_0 \ {}^t\boldsymbol{a}_1 \ {}^t\boldsymbol{a}_2) \in SO_0(1, 2)$, where

$$SO_0(1,2) = \left\{ A = \begin{pmatrix} a_0^0 & a_0^1 & a_0^2 \\ a_1^0 & a_1^1 & a_1^2 \\ a_2^0 & a_2^1 & a_2^2 \end{pmatrix} \mid {}^t A I_{1,2} A = I_{1,2}, \ a_0^0 \ge 1 \right\}$$

is the *Lorentz group*, where

$$I_{1,2} = \begin{pmatrix} -1 & 0 & 0\\ 0 & 1 & 0\\ 0 & 0 & 1 \end{pmatrix}.$$

For any $A = ({}^{t}\boldsymbol{a}_{0} {}^{t}\boldsymbol{a}_{1} {}^{t}\boldsymbol{a}_{2}) \in SO_{0}(1,2), \{\boldsymbol{a}_{0},\boldsymbol{a}_{1},\boldsymbol{a}_{2}\}$ is a pseudo orthonormal basis of \mathbb{R}^{3}_{1} . Then $\boldsymbol{\ell} = \boldsymbol{a}_{0} + \boldsymbol{a}_{2}$ is lightlike. It follows that we have $HC(\boldsymbol{\ell},-1) = HC(\boldsymbol{a}_{0} + \boldsymbol{a}_{2},-1)$ such that $\boldsymbol{a}_{0} \in HC(\boldsymbol{a}_{0} + \boldsymbol{a}_{2},-1)$ and \boldsymbol{a}_{1} is tangent to $HC(\boldsymbol{a}_{0} + \boldsymbol{a}_{2},-1)$ at \boldsymbol{a}_{0} . Moreover, we have $\boldsymbol{a}_{0} \in HL(\boldsymbol{a}_{2})$ and \boldsymbol{a}_{1} is tangent to $HL(\boldsymbol{a}_{2})$ at \boldsymbol{a}_{0} . Then we have the following lemma.

Lemma 3.1. With the above notation, we have (1) $HC(\ell, -1) = \left\{ x = a_0 + ra_1 + \frac{1}{2}r^2(a_0 + a_2) \mid r \in \mathbb{R} \right\}.$ (2) $HL(a_2) = \{ \sqrt{r^2 + 1}a_0 + ra_1 \mid r \in \mathbb{R} \}.$

Proof. (1) For any $\boldsymbol{x} \in HC(\boldsymbol{\ell}, -1)$, there exist $\alpha, \beta, \gamma \in \mathbb{R}$ such that

$$\boldsymbol{x} = \alpha \boldsymbol{a}_0 + \beta \boldsymbol{a}_1 + \gamma \boldsymbol{a}_2 \ (\alpha \ge 1).$$

We put $\beta = r$. Since $\langle \boldsymbol{x}, \boldsymbol{\ell} \rangle = -\alpha + \gamma = -1$, we have $\alpha = \gamma + 1$. Moreover, we also have $\langle \boldsymbol{x}, \boldsymbol{x} \rangle = -\alpha^2 + \beta^2 + \gamma^2 = -(\gamma + 1)^2 + r^2 + \gamma^2 = -1$, so that $\gamma = \frac{1}{2}r^2$. Thus,

$$x = a_0 + ra_1 + \frac{1}{2}r^2(a_0 + a_2)$$

holds. For the converse, we can easily show that $\langle \boldsymbol{x}, \boldsymbol{x} \rangle = -1$ and $\langle \boldsymbol{x}, \boldsymbol{\ell} \rangle = -1$ for the above vector.

(2) For any $\boldsymbol{x} \in HL(\boldsymbol{a}_2)$, there exist $\alpha, \beta, \gamma \in \mathbb{R}$ such that

$$\boldsymbol{x} = \alpha \boldsymbol{a}_0 + \beta \boldsymbol{a}_1 + \gamma \boldsymbol{a}_2 \ (\alpha \ge 1).$$

Since $\langle \boldsymbol{x}, \boldsymbol{a}_2 \rangle = 0$, $\gamma = 0$. If we put $\beta = r$, then we have $\langle \boldsymbol{x}, \boldsymbol{x} \rangle = -\alpha^2 + r^2 = -1$, so that $\alpha = \pm \sqrt{r^2 + 1}$. Since $\alpha \ge 1$, we have $\alpha = \sqrt{r^2 + 1}$. By a straightforward calculation, the converse holds.

It is known that a horocycle $\Psi(HC(\boldsymbol{a}_0 + \boldsymbol{a}_2, -1))$ in the Poincaré disc D is a Euclidean circle tangent to the ideal boundary $S^1 = \{\boldsymbol{x} \in \mathbb{R}^2_0 \mid ||\boldsymbol{x}|| = 1\}$. It is also known that a hyperbolic line $\Psi(HL(\boldsymbol{a}_2))$ is a Euclidean circle or a Euclidean line orthogonal to the ideal boundary (cf., [11]). By these reasons, a horocycle is called a *horizontal pseudo-line* and a hyperbolic-line is called an *orthogonal pseudo-line* respectively. We now define a ϕ -slant pseudo-line by

$$SL(\boldsymbol{n}_{\phi}, -\cos\phi) = \left\{ \boldsymbol{x} \in H^2_+(-1) \mid \langle \boldsymbol{x}, \boldsymbol{n}_{\phi} \rangle = -\cos\phi \right\},$$

where $\mathbf{n}_{\phi}(t) = \cos \phi \mathbf{a}_0 + \mathbf{a}_2$, $\phi \in [0, \pi/2]$. Since $\langle \mathbf{n}_{\phi}, \mathbf{n}_{\phi} \rangle = \sin^2 \phi > 0$, \mathbf{n}_{ϕ} is spacelike. Thus, $SL(\mathbf{n}_{\phi}, -\cos \phi) = H^2_+(-1) \cap P(\mathbf{n}_{\phi}, -\cos \phi)$ is an equidistant curve. Moreover, $\mathbf{a}_0 \in SL(\mathbf{n}_{\phi}, -\cos \phi)$ and \mathbf{a}_1 is tangent to $SL(\mathbf{n}_{\phi}, -\cos \phi)$ at \mathbf{a}_0 . Then $SL(\mathbf{n}_{\pi/2}, -\cos(\pi/2)) = HL(\mathbf{a}_2)$ and $SL(\mathbf{n}_0, -\cos 0) = HC(\mathbf{a}_0 + \mathbf{a}_2, -1)$. We have the following parametrization of a ϕ -slant pseudo-line.

Lemma 3.2. With the same notations as those in Lemma 3.1, we have

$$SL(\boldsymbol{n}_{\phi}, -\cos\phi) = \left\{ \boldsymbol{a}_{0} + r\boldsymbol{a}_{1} + \frac{\sqrt{r^{2}\sin^{2}\phi + 1} - 1}{\sin^{2}\phi} (\boldsymbol{a}_{0} + \cos\phi\boldsymbol{a}_{2}) \mid r \in \mathbb{R} \right\}.$$

Proof. We consider a point $\boldsymbol{x} \in SL(\boldsymbol{n}_{\phi}(t), -\cos \phi)$. Since $\{\boldsymbol{a}_0, \boldsymbol{a}_1, \boldsymbol{a}_2\}$ is a pseudo-orthonormal basis of \mathbb{R}^3_1 , There exist $\alpha, \beta, \gamma \in \mathbb{R}$ such that $\boldsymbol{x} = \alpha \boldsymbol{a}_0 + \beta \boldsymbol{a}_1 + \gamma \boldsymbol{a}_2$, $(\alpha \geq 1)$. Therefore, we have

$$\langle \boldsymbol{x}, \boldsymbol{n}_{\phi} \rangle = \langle \alpha \boldsymbol{a}_{0} + \beta \boldsymbol{a}_{1} + \gamma \boldsymbol{a}_{2}, \cos \phi \boldsymbol{a}_{0} + \boldsymbol{a}_{2} \rangle$$
$$= -\cos \phi \alpha + \gamma = -\cos \phi$$

Thus, we have $\gamma = \cos \phi(\alpha - 1)$. Moreover, $\langle \boldsymbol{x}, \boldsymbol{x} \rangle = -\alpha^2 + \beta^2 + \gamma^2 = -\alpha^2 + \beta^2 + \cos^2 \phi(\alpha - 1)^2 = -1$. It follows that

$$\alpha = \frac{1}{\sin^2 \phi} (\pm \sqrt{\beta^2 \sin^2 \phi + 1} - \cos^2 \phi).$$

If we choose $\alpha = -\frac{1}{\sin^2 \phi} (\sqrt{\beta^2 \sin^2 \phi + 1} + \cos^2 \phi)$, then $\alpha < 0$. It contradicts to $\alpha \ge 1$. Hence, we have

$$\alpha = \frac{1}{\sin^2 \phi} (\sqrt{\beta^2 \sin^2 \phi + 1} - \cos^2 \phi), \ \gamma = \frac{\cos \phi}{\sin^2 \phi} (\sqrt{\beta^2 \sin^2 \phi + 1} - 1).$$

We put $\beta = r$. Then

$$\begin{aligned} \boldsymbol{x} &= \frac{1}{\sin^2 \phi} (\sqrt{r^2 \sin^2 \phi + 1} - \cos^2 \phi) \boldsymbol{a}_0 + r \boldsymbol{a}_1 + \frac{\cos \phi}{\sin^2 \phi} (\sqrt{r^2 \sin^2 \phi + 1} - 1) \boldsymbol{a}_2 \\ &= \boldsymbol{a}_0 + r \boldsymbol{a}_1 + \frac{1}{\sin^2 \phi} (\sqrt{r^2 \sin^2 \phi + 1} - 1) (\boldsymbol{a}_0 + \cos \phi \boldsymbol{a}_2) \end{aligned}$$

For the converse, we have $\langle \boldsymbol{x}, \boldsymbol{x} \rangle = -1$, $\langle \boldsymbol{x}, \boldsymbol{n}_{\phi} \rangle = -\cos \phi$ and $(\sqrt{r^2 \sin^2 \phi + 1} - \cos^2 \phi) / \sin^2 \phi \ge 1$. Then $\boldsymbol{x} \in SL(\boldsymbol{n}_{\phi}(t), -\cos \phi)$.

Remark 3.3. We can show $\lim_{\phi\to 0} (\sqrt{r^2 \sin^2 \phi + 1} - 1) / \sin^2 \phi = r^2/2$. In [10] the third author showed that the angle between $\Psi(SL(\boldsymbol{n}_{\phi}))$ and the ideal boundary S^1 of the Poincaré disc D at an intersection point is equal to ϕ . This is the reason why we call $SL(\boldsymbol{n}_{\phi})$ the ϕ -slant pseudo line.

4 One-parameter families of pseudo-lines

In this section we consider one-parameter families of pseudo-lines. By Lemmas 3.1 and 3.2, we consider a one-parameter family of pseudo-orthonormal bases of \mathbb{R}^3_1 . Let $A : J \longrightarrow SO_0(1,2)$ be a C^{∞} -mapping. If we write $A(t) = ({}^t \boldsymbol{a}_0(t) \; {}^t \boldsymbol{a}_1(t) \; {}^t \boldsymbol{a}_2(t))$, then $\{\boldsymbol{a}_0(t), \boldsymbol{a}_1(t), \boldsymbol{a}_2(t)\}$ is a one-parameter family of pseudo-orthonormal bases of \mathbb{R}^3_1 . We call it a *pseudo-orthonormal moving frame*

of \mathbb{R}^3_1 . By the standard arguments, we can show the following Frenet-Serret type formulae for the pseudo-orthonormal moving frame $\{a_0(t), a_1(t), a_2(t)\}$:

$$\begin{pmatrix} \mathbf{a}_0'(t) \\ \mathbf{a}_1'(t) \\ \mathbf{a}_2'(t) \end{pmatrix} = \begin{pmatrix} 0 & c_1(t) & c_2(t) \\ c_1(t) & 0 & c_3(t) \\ c_2(t) & -c_3(t) & 0 \end{pmatrix} \begin{pmatrix} \mathbf{a}_0(t) \\ \mathbf{a}_1(t) \\ \mathbf{a}_2(t) \end{pmatrix}$$

Here,

$$\begin{cases} c_1(t) = \langle \boldsymbol{a}'_0(t), \boldsymbol{a}_1(t) \rangle \\ c_2(t) = \langle \boldsymbol{a}'_0(t), \boldsymbol{a}_2(t) \rangle \\ c_3(t) = \langle \boldsymbol{a}'_1(t), \boldsymbol{a}_2(t) \rangle \end{cases}$$

Then, the matrix $C(t) = \begin{pmatrix} 0 & c_1(t) & c_2(t) \\ c_1(t) & 0 & c_3(t) \\ c_2(t) & -c_3(t) & 0 \end{pmatrix}$ is an element of Lie algebra

 $\mathfrak{so}(1,2)$ of the Lorentz group $SO_0(1,2)$. The above Frenet-Serret type formulae are written by $A'(t)A^{-1}(t) = C(t)$. For any C^{∞} -mapping $C : J \longrightarrow \mathfrak{so}(1,2)$ and $A_0 \in SO_0(1,2)$, we can apply the unique existence theorem for systems of linear ordinary differential equations, so that there exists a unique $A(t) \in SO_0(1,2)$ such that $A(0) = A_0$ and $A'(t)A^{-1}(t) = C(t)$.

We now consider a mapping $g_{\phi}: I \times J \longrightarrow H^2_+$, where $g_{\phi}(r, t)$ is defined by

$$\begin{cases} \boldsymbol{a}_0(t) + r\boldsymbol{a}_1(t) + \frac{\sqrt{r^2 \sin^2 \phi + 1} - 1}{\sin^2 \phi} (\boldsymbol{a}_0(t) + \cos \phi \boldsymbol{a}_2(t)) & \text{if } \phi \neq 0, \\ \boldsymbol{a}_0(t) + r\boldsymbol{a}_1(t) + \frac{r^2}{2} (\boldsymbol{a}_0(t) + \boldsymbol{a}_2(t)) & \text{if } \phi = 0, \end{cases}$$

where $I, J \subset \mathbb{R}$ are intervals. Then we have $SL(\mathbf{n}_{\phi}(t), -\cos \phi) = \{g_{\phi}(I \times \{t\}) \mid t \in J\}$, for $\mathbf{n}_{\phi}(t) = \cos \phi \mathbf{a}_0(t) + \mathbf{a}_2(t)$. Thus g_{ϕ} is a one-parameter family of ϕ -slant pseudo-lines. Moreover, g_0 is a one-parameter family of horocycles and $g_{\pi/2}$ is a one-parameter family of hyperbolic lines.

5 Height functions

For a one parameter family of ϕ -slant pseudo-lines g_{ϕ} , we define a *family of* height functions $H : J \times H^2_+ \longrightarrow \mathbb{R}$ by $H(t, \boldsymbol{x}) = \langle \boldsymbol{x}, \boldsymbol{n}_{\phi}(t) \rangle + \cos \phi$. Then we have the following proposition.

Proposition 5.1. For $g_{\phi}: I \times J \longrightarrow H^2_+$, we have the following: (1) $H(t, \boldsymbol{x}) = 0$ if and only if there exists $r \in I$ such that $\boldsymbol{x} = g_{\phi}(r, t)$, (2) $H(t, \boldsymbol{x}) = \frac{\partial H}{\partial t}(t, \boldsymbol{x}) = 0$ if and only if there exists $r \in I$ such that $\boldsymbol{x} = g_{\phi}(r, t)$ and $-\sqrt{r^2 \sin^2 \phi + 1}c_2(t) + r(\cos \phi c_1(t) - c_3(t)) = 0.$ *Proof.* (1) If $H(t, \boldsymbol{x}) = 0$, then $\langle \boldsymbol{x}, \boldsymbol{n}_{\phi}(t) \rangle = -\cos \phi$, $\boldsymbol{x} \in H^2_+(-1)$. Thus, there exists $r \in I$ such that $\boldsymbol{x} = g_{\phi}(r, t)$. The converse also holds.

(2)Since $H(t, \boldsymbol{x}) = 0$, there exists $r \in I$ such that $\boldsymbol{x} = g_{\phi}(r, t)$. Suppose that $\phi \neq 0$. Since $\boldsymbol{n}_{\phi}'(t) = c_2(t)\boldsymbol{a}_0(t) + (\cos\phi c_1(t) - c_3(t))\boldsymbol{a}_1(t) + \cos\phi c_2(t)\boldsymbol{a}_2(t)$,

$$\begin{aligned} \frac{\partial H}{\partial t}(t, \boldsymbol{x}) &= \langle \boldsymbol{x}, \boldsymbol{n}_{\phi}'(t) \rangle \\ &= \left\langle \frac{\sqrt{r^2 \sin^2 \phi + 1} - \cos^2 \phi}{\sin^2 \phi} \boldsymbol{a}_0(t) + r \boldsymbol{a}_1(t) \right. \\ &+ \frac{\cos \phi (\sqrt{r^2 \sin^2 \phi + 1} - 1)}{\sin^2 \phi} \boldsymbol{a}_2(t), \ \boldsymbol{n}_{\phi}'(t) \right\rangle \\ &= -\sqrt{r^2 \sin^2 \phi + 1} c_2(t) + r \left(\cos \phi c_1(t) - c_3(t) \right). \end{aligned}$$

If $\phi = 0$, then

$$\frac{\partial H}{\partial t}(t, \boldsymbol{x}) = \langle \boldsymbol{x}, \boldsymbol{n}'_0(t) \rangle = -c_2(t) + r(c_1(t) - c_3(t)).$$

This completes the proof.

We now review some general results on the singularity theory for families of function germs. Detailed descriptions are found in the book[2]. Let $F : (\mathbb{R} \times \mathbb{R}^r, (s_0, x_0)) \to \mathbb{R}$ be a function germ. We call F an r-parameter unfolding of f, where $f(s) = F_{x_0}(s, x_0)$. We say that f has an A_k -singularity at s_0 if $f^{(p)}(s_0) =$ 0 for all $1 \leq p \leq k$, and $f^{(k+1)}(s_0) \neq 0$. We also say that f has an $A_{\geq k}$ singularity at s_0 if $f^{(p)}(s_0) = 0$ for all $1 \leq p \leq k$. Let F be an unfolding of f and f(s) has an A_k -singularity $(k \geq 1)$ at s_0 . We denote the (k-1)-jet of the partial derivative $\frac{\partial F}{\partial x_i}$ at s_0 by $j^{(k-1)}(\frac{\partial F}{\partial x_i}(s, x_0))(s_0) = \sum_{j=0}^{k-1} \alpha_{ji}(s-s_0)^j$ for $i = 1, \ldots, r$. Then F is called an \mathcal{R} -versal unfolding if the $k \times r$ matrix of coefficients $(\alpha_{ji})_{j=0,\ldots,k-1;i=1,\ldots,r}$ has rank k $(k \leq r)$. We introduce an important set concerning the unfoldings relative to the above notions. The discriminant set of F is the set

$$\mathcal{D}_F = \left\{ x \in \mathbb{R}^r | \text{there exists } s \text{ such that } F(s, x) = \frac{\partial F}{\partial s}(s, x) = 0 \right\}.$$

Then we have the following classification (cf., [2]).

Theorem 5.2. Let $F : (\mathbb{R} \times \mathbb{R}^r, (s_0, x_0)) \to \mathbb{R}$ be an *r*-parameter unfolding of f(s) which has an A_k singularity at s_0 (k = 1, 2). Suppose that F is an \mathcal{R} -versal unfolding.

(1) If k = 1, then \mathcal{D}_F is locally diffeomorphic to \mathbb{R}^{r-1} .

(2) If k = 2, then \mathcal{D}_F is locally diffeomorphic to $C \times \mathbb{R}^{r-2}$. Here, $C = \{(x_1, x_2) \in (\mathbb{R}^2, 0) \mid x_1 = t^2, x_2 = t^3, t \in (\mathbb{R}, 0) \}$ is the ordinary cusp.

By Proposition 5.1, the discriminant set \mathcal{D}_H of H is

$$D_H = \left\{ g_\phi(r,t) \mid -\sqrt{r^2 \sin^2 \phi + 1} c_2(t) + r \left(\cos \phi c_1(t) - c_3(t) \right) = 0 \right\}$$

Suppose $c_2(t) \neq 0, \cos \phi c_1(t) - c_3(t) \neq 0$. If

$$-\sqrt{r^2 \sin^2 \phi + 1}c_2(t) + r\left(\cos \phi c_1(t) - c_3(t)\right) = 0,$$

then we have

$$r = \pm \frac{c_2(t)}{\sqrt{(\cos\phi c_1(t) - c_3(t))^2 - (\sin\phi c_2(t))^2}}.$$

If $r = -c_2(t)/\sqrt{(\cos\phi c_1(t) - c_3(t))^2 - (\sin\phi c_2(t))^2}$, then
 $-\sqrt{r^2 \sin^2\phi + 1}c_2(t) + r(\cos\phi c_1(t) - c_3(t)) \neq 0$,

so that

$$r = \frac{c_2(t)}{\sqrt{(\cos\phi c_1(t) - c_3(t))^2 - (\sin\phi c_2(t))^2}}.$$

For $\phi = 0$, we can also choose $r = c_2(t)/(c_1(t) - c_3(t))$. Therefore, if $\phi \neq 0$, then

$$D_H = \left\{ g_{\phi}(r,t) \mid r = \frac{c_2(t)}{\sqrt{(\cos\phi c_1(t) - c_3(t))^2 - (\sin\phi c_2(t))^2}}, c_2(t) \neq 0, \\ \cos\phi c_1(t) - c_3(t) \neq 0 \right\}.$$

Under the assumptions that $c_2(t) \neq 0$ and $\cos \phi c_1(t) - c_3(t) \neq 0$, we have a $g[\phi]: J \longrightarrow H^2_+$, where $g[\phi](t)$ is defined by

$$\begin{cases} \boldsymbol{a}_0(t) + r(t)\boldsymbol{a}_1(t) + \frac{\sqrt{r(t)^2 \sin^2 \phi + 1} - 1}{\sin^2 \phi} (\boldsymbol{a}_0(t) + \cos \phi \boldsymbol{a}_2(t)) & \text{if } \phi \neq 0, \\ \boldsymbol{a}_0(t) + r(t)\boldsymbol{a}_1(t) + \frac{r(t)^2}{2} (\boldsymbol{a}_0(t) + \boldsymbol{a}_2(t)) & \text{if } \phi = 0. \end{cases}$$

Here

$$r(t) = \begin{cases} \frac{c_2(t)}{\sqrt{(\cos\phi c_1(t) - c_3(t))^2 - (\sin\phi c_2(t))^2}} & \text{if } \phi \neq 0, \\ \frac{c_2(t)}{c_1(t) - c_3(t)} & \text{if } \phi = 0. \end{cases}$$

Then $g[\phi](t)$ is a parametrization of D_H and it is the envelope of the family of ϕ -slant pseudo lines $\{SL(\mathbf{n}_{\phi}(t), -\cos \phi)\}_{t \in J}$.

In order to classify the singularities of $g[\phi]$, we apply the theory of unfoldings to H. For any $(r_0, t_0) \in I \times J$, we put $\mathbf{x}_0 = g_{\phi}(r_0, t_0)$ and consider the function germ $h_{\mathbf{x}_0}$: $(J, t_0) \longrightarrow (\mathbb{R}, 0)$ defined by

$$h_{\boldsymbol{x}_0}(t_0) = H(t_0, \boldsymbol{x}_0) = \langle \boldsymbol{x}_0, \boldsymbol{n}_{\boldsymbol{\phi}}(t_0) \rangle + \cos \phi$$

Then the germ of H at (t_0, \boldsymbol{x}_0) is a two-dimensional unfolding of $h_{\boldsymbol{x}_0}$. We now try to search the conditions for $h_{\boldsymbol{x}_0}$ has the A_k -singularity, (k = 1, 2). If $\phi \neq 0$, then we define two invariants

$$\begin{cases} \delta[\phi]_{1}(t) = -\sqrt{r^{2} \sin^{2} \phi + 1}c'_{2} + r(\cos \phi c'_{1} - c'_{3}) + rc_{2}(c_{1} - \cos \phi c_{3}) \\ -c_{1}(\cos \phi c_{1} - c_{3}) - \cos \phi c_{2}^{2} \\ + \frac{\sqrt{r^{2} \sin^{2} \phi + 1} - 1}{\sin^{2} \phi} (\cos \phi c_{1} - c_{3})(\cos \phi c_{3} - c_{1}), \\ \delta[\phi]_{2}(t) = -\sqrt{r^{2} \sin^{2} \phi + 1}c''_{2} \\ + r(\cos \phi c''_{1} - c''_{3}) + r\left(c_{2}(c'_{1} - \cos \phi c'_{3}) + 2c'_{2}(c_{1} - \cos \phi c_{3})\right) \\ -2c_{1}(\cos \phi c'_{1} - c'_{3}) - c'_{1}(\cos \phi c_{1} - c_{3}) - 3\cos \phi c_{2}c'_{2} \\ + \frac{\sqrt{r^{2} \sin^{2} \phi + 1} - 1}{\sin^{2} \phi} \left((\cos \phi c_{1} - c_{3})(\cos \phi c'_{3} - c'_{1}) \\ + 2(\cos \phi c'_{1} - c'_{3})(\cos \phi c_{3} - c_{1})\right), \end{cases}$$

where $c_i = c_i(t)$, (i = 1, 2, 3) and r = r(t). For $\phi = 0$, we also define

$$\begin{cases} \delta[0]_1(t) = -c'_2 + r(c'_1 - c'_3) - c_1(c_1 - c_3) - \frac{1}{2}c_2^2, \\ \delta[0]_2(t) = -c''_2 + r(c''_1 - c''_3) + rc_2(c'_1 - c'_3) - 2c_1(c'_1 - c'_3) \\ -c'_1(c_1 - c_3) - c_2c'_2 + \frac{3}{2}rc_2(c'_3 - c'_1). \end{cases}$$

We remark that $\lim_{\phi\to 0} \delta[\phi]_1(t) = \delta[0]_1(t)$ and $\lim_{\phi\to 0} \delta[\phi]_2(t) = \delta[0]_2(t)$. We expect that $\delta[\phi]_1(t) = \delta[\phi]_2(t) = 0$ if and only if $\delta[\phi]_1(t) = \delta[\phi]'_1(t) = 0$. However, we can only show this relation for a special case (cf., §6).

Proposition 5.3. We have the following assertions:

(1) $h'_{\boldsymbol{x}_0}(t_0) = 0$ always holds, (2) $h''_{\boldsymbol{x}_0}(t_0) = 0$ if and only if $\delta[\phi]_1(t_0) = 0$, (3) $h''_{\boldsymbol{x}_0}(t_0) = h'''_{\boldsymbol{x}_0}(t_0) = 0$ if and only if $\delta[\phi]_1(t_0) = \delta[\phi]_2(t_0) = 0$. *Proof.* Assertion (1) holds by Proposition 5.1, (2).

$$(2)$$
 Since

$$\begin{aligned} \boldsymbol{n}_{\phi}^{\prime\prime} &= (\boldsymbol{n}_{\phi}^{\prime})^{\prime} \\ &= c_{2}^{\prime} \boldsymbol{a}_{0} + (\cos \phi c_{1}^{\prime} - c_{3}^{\prime}) \boldsymbol{a}_{1} + \cos \phi c_{2}^{\prime} \boldsymbol{a}_{2} + (c_{1}(\cos \phi c_{1} - c_{3}) + \cos \phi c_{2}^{2}) \boldsymbol{a}_{0} \\ &+ c_{2}(c_{1} - \cos \phi c_{3}) \boldsymbol{a}_{1} + (c_{3}(\cos \phi c_{1} - c_{3}) + c_{2}^{2}) \boldsymbol{a}_{2}, \end{aligned}$$

we have

$$\begin{split} \frac{\partial^2 H}{\partial t^2}(t, \boldsymbol{x}) &= \langle \boldsymbol{x}, \boldsymbol{n}_{\phi}' \rangle \\ &= -\sqrt{s^2 \sin^2 \phi + 1} c_2' + s(\cos \phi c_1' - c_3') \\ &+ \frac{(\cos^2 \phi - \sqrt{s^2 \sin^2 \phi + 1}) \left(c_1(\cos \phi c_1 - c_3) + \cos \phi c_2^2\right)}{\sin^2 \phi} \\ &+ sc_2(c_1 - \cos \phi c_3) + \frac{\cos \phi (\sqrt{s^2 \sin^2 \phi + 1} - 1) \left(c_3(\cos \phi c_1 - c_3) + c_2^2\right)}{\sin^2 \phi} \\ &= -\sqrt{s^2 \sin^2 \phi + 1} c_2' + s(\cos \phi c_1' - c_3') - c_1(\cos \phi c_1 - c_3) - \cos \phi c_2^2 \\ &+ \frac{(1 - \sqrt{s^2 \sin^2 \phi + 1}) \left(c_1(\cos \phi c_1 - c_3) + \cos \phi c_2^2\right)}{\sin^2 \phi} \\ &+ \frac{\cos \phi (\sqrt{s^2 \sin^2 \phi + 1} - 1) \left(c_3(\cos \phi c_1 - c_3) + c_2^2\right)}{\sin^2 \phi} \\ &= -\sqrt{s^2 \sin^2 \phi + 1} c_2' + s(\cos \phi c_1' - c_3') + sc_2(c_1 - \cos \phi c_3) \\ &- c_1(\cos \phi c_1 - c_3) - \cos \phi c_2^2 + \frac{(\sqrt{s^2 \sin^2 \phi + 1} - 1)(\cos \phi c_1 - c_3)(\cos \phi c_3 - c_1)}{\sin^2 \phi} \\ &= \delta[\phi]_1(t). \end{split}$$

(3) We have

$$\boldsymbol{n}_{\phi}^{\prime\prime\prime} = c_2^{\prime\prime} \boldsymbol{a}_0 + (\cos\phi c_1^{\prime\prime} - c_3^{\prime\prime}) \boldsymbol{a}_1 + \cos\phi c_2^{\prime\prime} \boldsymbol{a}_2 + 2 \left(c_2^{\prime} \boldsymbol{a}_0^{\prime} + (\cos\phi c_1^{\prime} - c_3^{\prime}) \boldsymbol{a}_1^{\prime} + \cos\phi c_2^{\prime} \boldsymbol{a}_2^{\prime} \right) + c_2 \boldsymbol{a}_0^{\prime\prime} + (\cos\phi c_1 - c_3) \boldsymbol{a}_1^{\prime\prime} + \cos\phi c_2 \boldsymbol{a}_2^{\prime\prime}.$$

Here, $a'_0 = c_1 a_1 + c_2 a_2$, $a'_1 = c_1 a_0 + c_3 a_2$, $a'_2 = c_2 a_0 - c_3 a_1$. Then

$$\begin{cases} \boldsymbol{a}_0'' = (c_1^2 + c_2^2)\boldsymbol{a}_0 + (c_1' - c_2c_3)\boldsymbol{a}_1 + (c_2' + c_1c_3)\boldsymbol{a}_2, \\ \boldsymbol{a}_1'' = (c_1' + c_2c_3)\boldsymbol{a}_0 + (c_1^2 - c_3^2)\boldsymbol{a}_1 + (c_3' + c_1c_2)\boldsymbol{a}_2, \\ \boldsymbol{a}_2'' = (c_2' - c_1c_3)\boldsymbol{a}_0 + (c_1c_2 - c_3')\boldsymbol{a}_1 + (c_2^2 - c_3^2)\boldsymbol{a}_2. \end{cases}$$

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Therefore,

$$\begin{aligned} \boldsymbol{n}_{\phi}^{\prime\prime\prime} &= c_2^{\prime\prime} \boldsymbol{a}_0 + (\cos \phi c_1^{\prime\prime} - c_3^{\prime\prime}) \boldsymbol{a}_1 + \cos \phi c_2^{\prime\prime} \boldsymbol{a}_2 \\ &+ \left(c_2 (c_1^2 + c_2^2 - c_3^2) + 2c_1 (\cos \phi c_1^{\prime} - c_3^{\prime}) + c_1^{\prime} (\cos \phi c_1 - c_3) + 3\cos \phi c_2 c_2^{\prime} \right) \boldsymbol{a}_0 \\ &+ \left((\cos \phi c_1 - c_3) (c_1^2 + c_2^2 - c_3^2) + c_2 (c_1^{\prime} - \cos \phi c_3^{\prime}) + 2c_2^{\prime} (c_1 - \cos \phi c_3) \right) \boldsymbol{a}_1 \\ &+ \left(\cos \phi c_2 (c_1^2 + c_2^2 - c_3^2) + 2c_3 (\cos \phi c_1^{\prime} - c_3^{\prime}) + c_3^{\prime} (\cos \phi c_1 - c_3) + 3c_2 c_2^{\prime} \right) \boldsymbol{a}_2. \end{aligned}$$

By the calculation similar to case (2), we have

$$\frac{\partial^3 H}{\partial t^3}(t, \boldsymbol{x}) = \delta[\phi]_2(t)$$

For the case $\phi = 0$, we also have the similar arguments to the above case. This completes the proof.

We have the following corollary.

Corollary 5.4. For h_{x_0} as the above proposition, we have the following: (1) h_{x_0} has the A_1 -singularity at $t = t_0$ if and only if $\delta[\phi]_1(t_0) \neq 0$.

(2) h_{x_0} has the A_2 -singularity at $t = t_0$ if and only if $\delta[\phi]_1(t_0) = 0$, $\delta[\phi]_2(t_0) \neq 0$. Then we have the following proposition.

Proposition 5.5. For $h_{\boldsymbol{x}_0}$ as the above proposition, we have the following: (1) If $h_{\boldsymbol{x}_0}$ has the A_1 -singularity, then H is a \mathcal{R} -versal unfolding of $h_{\boldsymbol{x}_0}$, (2) If $h_{\boldsymbol{x}_0}$ has the A_2 -singularity, then H is a \mathcal{R} -versal unfolding of $h_{\boldsymbol{x}_0}$. *Proof.* We consider a parametrization of H^2_+ defined by

$$\psi(x_2, x_3) = (\sqrt{x_2^2 + x_3^2 + 1}, x_2, x_3).$$

Then we have

$$H(t, x_2, x_3) = H(t, \psi(x_2, x_3)) = \langle \psi(x_2, x_3), \mathbf{n}_{\phi}(t) \rangle + \cos \phi.$$

We write $\boldsymbol{n}_{\phi}(t) = (n_{\phi 1}(t), n_{\phi 2}(t), n_{\phi 3}(t))$ and have

$$\frac{\partial \dot{H}}{\partial x_i}(t, x_2, x_3) = n_{\phi i}(t) - \frac{x_i}{\sqrt{x_2^2 + x_3^2 + 1}} n_{\phi 1}(t) \quad (i = 2, 3).$$

Moreover, we have

$$\frac{\partial}{\partial t}\frac{\partial H}{\partial x_i}(t, x_2, x_3) = n'_{\phi i}(t) - \frac{x_i}{\sqrt{x_2^2 + x_3^2 + 1}}n'_{\phi 1}(t).$$

We write that $x_0 = (x_{01}, x_{02}, x_{03})$. Then the 1-jet of $(\partial \tilde{H} / \partial x_i)(t, x_{02}, x_{03})$ at $t = t_0$ is

$$\frac{\partial \tilde{H}}{\partial x_i}(t, x_{02}, x_{03}) = \frac{\partial \tilde{H}}{\partial x_i}(t_0, x_{02}, x_{03}) + \frac{1}{2}\frac{\partial}{\partial t}\frac{\partial \tilde{H}}{\partial x_i}(t_0, x_{02}, x_{03})(t - t_0).$$

From now on, we remove (t_0) for abbreviation.

(1) Since h_{x_0} has the A₁-singularity, we show that the rank of the matrix

$$\left(n_{\phi 2} - \frac{x_{02}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1} \quad n_{\phi 3} - \frac{x_{03}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1}\right)$$

Is equal to one. If the rank is zero, then

$$n_{\phi 2} - \frac{x_{02}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1} = -\frac{\cos \phi x_{02}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} = 0$$
$$n_{\phi 3} - \frac{x_{03}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1} = 1 - \frac{\cos \phi x_{03}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} = 0$$

Thus, we have the sum of the power of the both equations

$$0 = \frac{\cos^2 \phi(x_{02}^2 + x_{03}^2) - (x_{02}^2 + x_{03}^2 + 1)}{x_{02}^2 + x_{03}^2 + 1}$$
$$= -\frac{\sin^2 \phi(x_{02}^2 + x_{03}^2) + 1}{x_{02}^2 + x_{03}^2 + 1} \neq 0.$$

This is a contradiction.

(2) Since $h_{\boldsymbol{x_0}}$ has the A_2 -singularity, we show that the rank of the matrix

$$B = \begin{pmatrix} n_{\phi 2} - \frac{x_{02}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1} & n_{\phi 3} - \frac{x_{03}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1} \\ n_{\phi 2}' - \frac{x_{02}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1}' & n_{\phi 03}' - \frac{x_{03}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1}' \end{pmatrix}$$

is equal to two. Since $\mathbf{n}_{\phi}' = c_2 \mathbf{a}_0 + (\cos \phi c_1 - c_3) \mathbf{a}_1 + \cos \phi c_2 \mathbf{a}_2$,

$$\det B = \begin{vmatrix} n_{\phi 2} - \frac{x_{02}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1} & n_{\phi 3} - \frac{x_{03}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n_{\phi 1} \\ n'_{\phi 2} - \frac{x_{02}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n'_{\phi 1} & n'_{\phi 3} - \frac{x_{03}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} n'_{\phi 1} \end{vmatrix}$$
$$= \frac{1}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} |x_0, n_{\phi}, n'_{\phi}|$$
$$= -\frac{(\cos \phi c_1 - c_3)\sqrt{s^2 \sin^2 \phi + 1} + s \sin^2 \phi c_2}}{\sqrt{x_{02}^2 + x_{03}^2 + 1}} \neq 0$$

This means that the rank of B is two.

It follows from Theorem 5.2 and Proposition 5.5, we have shown the following theorem.

Theorem 5.6. Let $\{a_0(t), a_1(t), a_2(t)\}_{t \in J}$ be pseudo-orthonormal moving frame of \mathbb{R}^3_1 . Suppose $c_2(t) \neq 0$ and $\cos \phi c_1(t) - c_3(t) \neq 0$. Then we have the following:

(1) The envelope $g[\phi]$ of the family of ϕ -slant pseudo-lines $SL(\mathbf{n}_{\phi}, -\cos\phi)$ is regular at a point $t = t_0$ if and only if $\delta[\phi]_1(t_0) \neq 0$,

(2) The envelope $g[\phi]$ of the family of ϕ -slant pseudo-lines $SL(\mathbf{n}_{\phi}, -\cos\phi)$ at a point $t = t_0$ is locally diffeomorphic to the cusp C if and only if $\delta[\phi]_1(t_0) = 0$, $\delta[\phi]_2(t_0) \neq 0$.

6 Slant evolutes of hyperbolic plane curves

There is the notion of hyperbolic evolutes of hyperbolic plane curves [5]. Let $\gamma: J \longrightarrow H^2_+$ be a unit speed curve, where we use the parameter $s \in J$ instead of t. We call $\mathbf{t}(s) = \gamma'(s)$ a unit tangent vector of γ at s. Since $\langle \gamma(s), \gamma(s) \rangle = -1$ we have $\langle \gamma(s), \mathbf{t}(s) \rangle = 0$. We define $\mathbf{e}(s) = \gamma(s) \wedge \mathbf{t}(s)$, which is called a unit binormal vector of γ at $s \in J$. Then we have $\langle \mathbf{e}(s), \mathbf{e}(s) \rangle = \langle \gamma(s) \wedge \mathbf{t}(s), \gamma(s) \rangle \wedge \mathbf{t}(s), \gamma(s) \rangle \langle \mathbf{t}(s), \mathbf{t}(s) \rangle + \langle \gamma(s), \mathbf{t}(s) \rangle^2 = 1$. Therefore, we have a pseudo-orthonormal moving frame $\{\gamma(s), -\mathbf{e}(s), \mathbf{t}(s)\}$ of \mathbb{R}^3_1 , which is called a hyperbolic Sabban frame along γ .

$$a_0(s) = \gamma(s) , \ a_1(s) = -e(s) , \ a_2(s) = t(s)$$

Then we have the following Frenet-Serret type formulae:

$$\begin{cases} \boldsymbol{\gamma}'(s) = \boldsymbol{t}(s) \\ \boldsymbol{t}'(s) = \boldsymbol{\gamma}(s) + \kappa_g(s)\boldsymbol{e}(s) \\ \boldsymbol{e}'(s) = -\kappa_g(s)\boldsymbol{t}(s), \end{cases}$$

where $\kappa_g(s) = |\boldsymbol{\gamma}(s), \boldsymbol{\gamma}'(s), \boldsymbol{\gamma}''(s)|$ is called the *geodesic curvature* of $\boldsymbol{\gamma}$. Since $\boldsymbol{a}_0(s) = \boldsymbol{\gamma}(s)$, $\boldsymbol{a}_1(s) = -\boldsymbol{e}(s)$, $\boldsymbol{a}_2(s) = \boldsymbol{t}(s)$, we have $c_1(s) = 0, c_2(s) = 1$ and $c_3(s) = -\langle \boldsymbol{a}_1(s), \boldsymbol{a}_2'(s) \rangle = \langle \boldsymbol{e}(s), \boldsymbol{t}'(s) \rangle = \langle \boldsymbol{\gamma}(s) \wedge \boldsymbol{t}(s), \boldsymbol{t}'(s) \rangle = |\boldsymbol{\gamma}(s), \boldsymbol{t}(s), \boldsymbol{t}'(s)| = |\boldsymbol{\gamma}(s), \boldsymbol{\tau}'(s)| = |\boldsymbol{\gamma}(s), \boldsymbol{\gamma}'(s)| = \kappa_g(s)$. In this case, the family of ϕ -slant pseudo-lines $g_\phi: I \times J \longrightarrow H^2_+(-1)$ is

$$g_{\phi}(r,s) = \begin{cases} \gamma(s) - r \, \boldsymbol{e}(s) + \frac{\sqrt{r^2 \sin^2 \phi + 1} - 1}{\sin^2 \phi} (\gamma(s) + \cos \phi \boldsymbol{t}(s)) & \text{if } \phi \neq 0, \\ \gamma(s) - r \, \boldsymbol{e}(s) + \frac{r^2}{2} (\gamma(s) + \boldsymbol{t}(s)) & \text{if } \phi = 0. \end{cases}$$

Therefore, the envelope $g[\phi]: J \longrightarrow H^2_+$ of g_{ϕ} is

$$g[\phi](s) = \begin{cases} \gamma(s) - r(s) \, \boldsymbol{e}(s) + \frac{\sqrt{r(s)^2 \sin^2 \phi + 1} - 1}{\sin^2 \phi} (\gamma(s) + \cos \phi \boldsymbol{t}(s)) \text{ if } \phi \neq 0, \\ \gamma(s) + \frac{1}{\kappa_g(s)} \boldsymbol{e}(s) + \frac{1}{2\kappa_g^2(s)} (\gamma(s) + \boldsymbol{t}(s)) \text{ if } \phi = 0, \end{cases}$$

where

$$r(s) = \frac{1}{\sqrt{\kappa_g^2(s) - \sin^2 \phi}}.$$

We call $g[\pi/2]$ a hyperbolic evolute and g[0] a horocyclic evolute of γ , respectively. For $s_0 \in J$, we define

$$\begin{cases} \sigma[\phi]_1(s_0) = \kappa'_g(s_0) + \frac{\cos\phi(\kappa_g(s_0)^2 - \sin^2\phi)}{\sqrt{\kappa_g^2(s_0) - \sin^2\phi} - \kappa_g(s_0)}, \\ \sigma[\phi]_2(s_0) = \kappa''_g(s_0) + \cos\phi\kappa'_g(s_0) \left(\frac{\sqrt{\kappa_g^2(s_0) - \sin^2\phi} + 2\kappa_g(s_0)}{\sqrt{\kappa_g^2(s_0) - \sin^2\phi} - \kappa_g(s_0)}\right) \end{cases}$$

In this case, by a straightforward calculation, we can show that $\sigma[\phi]'_1(s) = \sigma[\phi]_2(s)$. Moreover, we can show that $\delta[\phi]_1(s) = \delta[\phi]_2(s) = 0$ if and only if $\sigma[\phi]_1(s) = \sigma[\phi]'_1(s) = 0$. As special cases, we have

$$\sigma[0]_1(s) = \kappa'_g(s) - \frac{1}{2}\kappa_g(s), \ \sigma[\pi/2]_1(s) = \kappa'_g(s).$$

As a corollary of Theorem 5.6, we have the following theorem.

Theorem 6.1. Let $\gamma : J \longrightarrow H^2_+(-1)$ be a unit speed curve with $\kappa_g(s)^2 - \sin^2 \phi > 0$. Then we have the following:

(1) $g[\phi]$ is a regular curve at $s = s_0$ if and only if $\sigma[\phi]_1(s_0) \neq 0$,

(2) $g[\phi]$ is locally diffeomorphic to the cusp C at $s = s_0$ if and only if

$$\sigma[\phi]_1(s_0) = 0 \text{ and } \sigma[\phi]'_1(s_0) \neq 0.$$

As a special case, we have the following corollary.

Corollary 6.2. Let $\gamma : J \longrightarrow H^2_+(-1)$ be a unit speed curve. (A) Suppose $\kappa_g^2 > 1$. Then we have the following (cf., [5]):

(1) The hyperbolic evolute $g[\pi/2]$ is a regular curve at $s = s_0$ if and only if $\kappa'_a(s) \neq 0$.

(2) The hyperbolic evolute $g[\pi/2]$ is locally diffeomorphic to the cusp C at $s = s_0$ if and only if

$$\kappa'_g(s_0) \neq 0 \text{ and } \kappa''_g(s_0) \neq 0.$$

(B) Suppose $\kappa_g \neq 0$. Then we have the following:

(1) The horocyclic evolute g[0] is a regular curve at $s = s_0$ if and only if $\kappa'_a(s) - \frac{1}{2}\kappa_g(s) \neq 0$.

(2) The horocyclic evolute g[0] is locally diffeomorphic to the cusp C at $s = s_0$ if and only if

$$\kappa'_g(s_0) - \frac{1}{2}\kappa_g(s_0) = 0 \text{ and } \kappa''_g(s_0) - \frac{1}{2}\kappa'_g(s_0) \neq 0.$$

The hyperbolic evolute is given by

$$g[\pi/2](s) = \begin{cases} \frac{-1}{\sqrt{\kappa_g^2(s) - 1}} (\kappa_g(s)\boldsymbol{\gamma}(s) + \boldsymbol{e}(s)) \text{ if } \kappa_g(s) < -1, \\ \frac{1}{\sqrt{\kappa_g^2(s) - 1}} (\kappa_g(s)\boldsymbol{\gamma}(s) - \boldsymbol{e}(s)) \text{ if } \kappa_g(s) > 1 \end{cases}$$

and the horocyclic evolute is

$$g[0](s) = \boldsymbol{\gamma}(s) + \frac{1}{\kappa_g(s)}\boldsymbol{e}(s) + \frac{1}{2\kappa_g^2(s)}(\boldsymbol{\gamma}(s) + \boldsymbol{t}(s)).$$

In [5] hyperbolic evolutes was introduced and the classified the singularities. Moreover, a de Sitter evolute of γ was introduced in [5], which is located in the de Sitter 2-space. It corresponds to points of $\gamma(s)$ with $\kappa_g^2(s) < 1$. Here we only consider families of hyperbolic lines, so that we do not consider de Sitter evolutes. It is also shown in [5] that $g[\pi/2](s)$ is a constant point if and only if γ is a part of a circle. This condition is also equivalent to $\kappa'_g(s) \equiv 0$. We have a natural question what is γ when g[0](s) is a constant point. Of course it is equivalent to

$$\kappa_g'(s) - \frac{1}{2}\kappa_g(s) \equiv 0.$$

The solution of the above differential equation is $\kappa_g(s) = ce^{s/2}$ for a constant real number c. The curvature tends to infinity, so that γ is a kind of spirals in $H^2_+(-1)$. If c = 1/2, the curve with the curvature $\frac{1}{2}c^{s/2}$ in the Euclidean plane is called a *Nielsen spiral*. So we call γ with $\kappa_g(s) = \frac{1}{2}c^{s/2}$ a hyperbolic Nielsen spiral. We have two open problems as follows:

(1) What is γ with $\sigma[\phi]_1(s) \equiv 0$?

(2) For a general one-parameter family of pseudo-lines, is it always true that $\delta[\phi]_1(t) = \delta[\phi]_2(t) = 0$ if and only if $\delta[\phi]_1(t) = \delta[\phi]'_1(t) = 0$?

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